

# Méthodes spectrales et preuves assistées par ordinateur pour l'étude d'EDP semi-linéaires sur $\mathbb{R}^m$

Jean-Philippe Lessard (McGill)

in collaboration with

Matthieu Cadiot (McGill)  
Jean-Christophe Nave (McGill)

SMAI Conference 2023, La Guadeloupe

# Computer-assisted proofs in dynamics

- Cesari [1964] Functional analysis and Galerkin's method
- Lanford [1982] Feigenbaum conjectures
- Fefferman & de la Llave [1986] Relativistic stability of matter
- Mischaikow & Mrozek [1995] Chaos in the Lorenz equations
- $\vdots$
- Jaquette & van den Berg [2019] Wright's conjecture (delay eq.)
- Wilczak & Zgliczyński [2020] Chaos in Kuramoto-Sivashinsky
- van den Berg, Breden, L. & van Veen [2021] Spontaneous periodic orbits in the Navier-Stokes equations
- Buckmaster, Cao-Labora & Gómez-Serrano [2022] Imploding solutions for 3D compressible fluids
- Chen & Hou [2022] Blowup in 3D Euler's equations

# What about PDEs defined on unbounded domains?

A motivation: develop a method to prove **constructively** existence of localized planar patterns, which arise in the study of

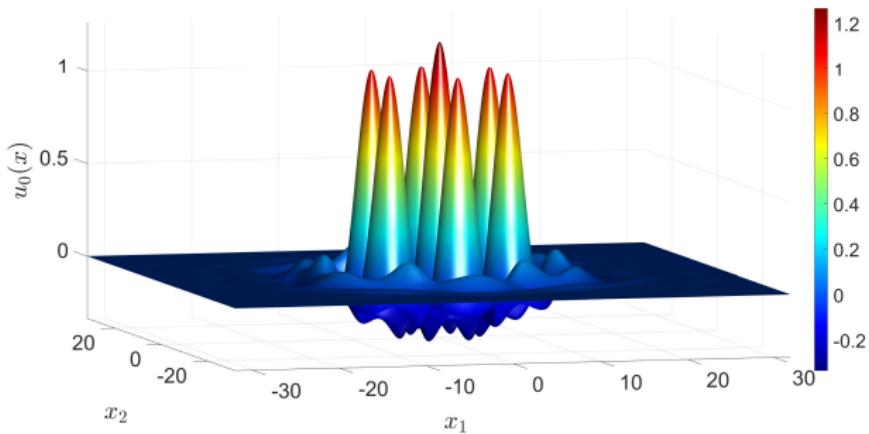
- dryland vegetation
- nonlinear optics
- phase-field crystals
- water waves
- neural field equations
- granular dynamics
- binary fluid convection & ferrofluids

## A mathematical challenge

Despite the prevalence and importance of localized planar patterns, little is known about them from a mathematical perspective.

# A motivating example on unbounded domains: non-radial localized stationary patterns in the Swift-Hohenberg PDE

$$(1 + \Delta)^2 u + \mu u + \nu_1 u^2 + \nu_2 u^3 = 0.$$



- D. Hill, J. Bramburger & D. Lloyd. *Approximate localised dihedral patterns near a Turing instability*. Nonlinearity, 36(5):2567, 2023.

# Formulation of the problem

We look for stationary planar solutions to

$$u_t = (1 + \Delta)^2 u + \mu u + \nu_1 u^2 + \nu_2 u^3 = 0$$

with  $u \rightarrow 0$  as  $|x| \rightarrow \infty$ . Assume that  $\mu > 0$ .

Let

$$\begin{aligned}\mathbb{L} &\stackrel{\text{def}}{=} \mu I_d + (I_d + \Delta)^2 \\ \mathbb{G}(u) &\stackrel{\text{def}}{=} \nu_1 u^2 + \nu_2 u^3 \\ \mathbb{F}(u) &\stackrel{\text{def}}{=} \mathbb{L}u + \mathbb{G}(u)\end{aligned}$$

then equivalently we have the following zero finding problem

$$\mathbb{F}(u) = 0 \quad \text{and} \quad u \rightarrow 0 \text{ as } |x| \rightarrow \infty.$$

# Main steps to get the computer-assisted proofs

- Develop a method (fully) based on Fourier series.
- Treat the problem on  $\mathbb{R}^2$  as a periodic problem (on a big enough square).
- Construct an approximate solution  $u_0$  on  $\mathbb{R}^2$  using Fourier series.
- Construct an approximate inverse for  $D\mathbb{F}(u_0)$  using Fourier series.
- Verify rigorously the hypotheses of a Newton-Kantorovich theorem close to  $u_0$  to show existence of an exact solution to  $\mathbb{F}(u) = 0$

# Operator representations

Let  $\Omega_0 \stackrel{\text{def}}{=} (-d, d)^2$ , the operators defined in the previous slides have a Fourier series representation on  $\Omega_0$ .

Operator	Fourier transform	Fourier series
$\mathbb{L} = \mu I_d + (I_d + \Delta)^2$	$l(\xi) = \mu + (1 -  \xi ^2)^2$	$L_p = \text{diag}\left(l\left(\frac{\pi}{d}n\right)\right)$
$u^2$	$\hat{u} * \hat{u}$	$U * U$
$\mathbb{F}(u) = \mathbb{L}u + \mathbb{G}(u)$	$l(\xi)\hat{u}(\xi) + \widehat{\mathbb{G}(u)}$	$F_p(U) = L_p U + G_p(U)$

# The Hilbert spaces $H^l$ and $X^l$

First ingredient: a natural Hilbert space for the solutions.

- **Assume**  $l(\xi) \stackrel{\text{def}}{=} \mathcal{F}(\mathbb{L})(\xi) = \mu + (1 - |\xi|^2)^2 > 0$
- Define  $H^l \stackrel{\text{def}}{=} \{u \in L^2, \|u\|_l < \infty\}$ , where

$$\|u\|_l^2 \stackrel{\text{def}}{=} \|\mathbb{L}u\|_2^2 = \int_{\mathbb{R}^2} l(\xi)^2 |\hat{u}(\xi)|^2 d\xi$$

- $H^l$  is a Banach algebra if  $\frac{1}{l} \in L^2 \implies \mathbb{G} : H^l \rightarrow L^2$  is well-defined
- Similarly,

$$X^l \stackrel{\text{def}}{=} \{U \in \ell^2, \|U\|_l < \infty\}$$

where  $\|U\|_l = \|L_p U\|_2$ .

# Elimination of the translation invariance

- There is a natural translation invariance of the set of solutions.
- Restrict to  $D_2$  symmetric functions to eliminate the invariance.
- $H_e \stackrel{\text{def}}{=} \{u \in H^l : u(x_1, x_2) = u(-x_1, x_2) = u(x_1, -x_2), \forall x_1, x_2 \in \mathbb{R}^2\}$
- $L_e$  is the restriction of  $L^2$  to  $D_2$  symmetric functions.

The zero finding problem becomes

$$\mathbb{F}(u) = \mathbb{L}u + \mathbb{G}(u) = 0 \quad \text{and} \quad u \in H_e$$

and  $\mathbb{F} : H_e \rightarrow L_e$ .

# Newton-Kantorovich approach

- Suppose we have an approximate solution  $u_0 \in H_e$
- Suppose that  $D\mathbb{F}(u_0)^{-1} : L_e \rightarrow H_e$  is a bounded linear operator

Prove that there exists  $r > 0$  such that  $\mathbb{T} : \overline{B_r(u_0)} \rightarrow \overline{B_r(u_0)}$

$$\mathbb{T}(u) \stackrel{\text{def}}{=} u - D\mathbb{F}(u_0)^{-1}\mathbb{F}(u),$$

is a contraction  $\implies$  there exists a unique solution in  $\overline{B_r(u_0)} \subset H_e$

# Newton-Kantorovich approach

## A Newton-Kantorovich type theorem

Consider  $u_0$  a numerical approximation of  $\mathbb{F}(u) = 0$  and suppose that  $\mathbb{A} : L_e \rightarrow H_e$  is an approximate inverse of  $D\mathbb{F}(u_0)$ . Consider bounds  $\mathcal{Y}_0$ ,  $\mathcal{Z}_1$  and  $\mathcal{Z}_2$  satisfying

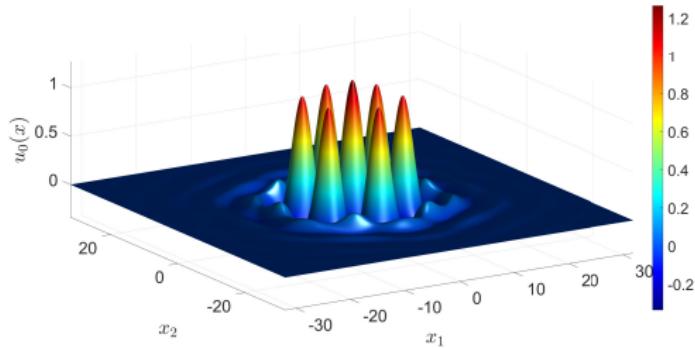
- $\|\mathbb{A}\mathbb{F}(u_0)\|_l \leq \mathcal{Y}_0$
- $\|I_d - \mathbb{A}D\mathbb{F}(u_0)\|_l \leq \mathcal{Z}_1$
- $\|\mathbb{A}(D\mathbb{F}(u_0) - D\mathbb{F}(v))\|_l \leq \mathcal{Z}_2(r)r$ , for all  $v \in \overline{B_r(u_0)}$ .

If there exists  $r > 0$  such that

$$\mathcal{Z}_1 < 1 \quad \text{and} \quad 2\mathcal{Y}_0\mathcal{Z}_2(r) < (1 - \mathcal{Z}_1)^2$$

then  $\mathbb{T} : \overline{B_r(u_0)} \rightarrow \overline{B_r(u_0)}$  is a contraction.

# Computation of the numerical approximation $u_0$



- $\text{supp}(u_0) \subset \Omega_0 \stackrel{\text{def}}{=} (-d, d)^2$ .
- Numerically, we can define  $u_0$  on  $\Omega_0$  by a Fourier series  $U_0 = (a_n)_{n \in \mathbb{Z}^2}$  where  $a_n = 0$  if  $|n|_\infty > N$  for some  $N \in \mathbb{N}$ .
- $u_0$  is fully determined by  $U_0$ :  $u_0(x) = \mathbf{1}_{\Omega_0}(x) \sum_{|n|_\infty \leq N} a_n e^{i \frac{\pi}{d} n \cdot x}$ .

# Computation of the numerical approximation $u_0$

- Assume that using a numerical scheme (e.g. Newton's method), we computed  $\tilde{U}_0$  such that  $F_p(\tilde{U}_0) \approx 0$  (Fourier problem on  $\Omega_0$ ).
- We build  $\Gamma$  such that  $\Gamma^*$  is full rank and such that the Fourier sequence

$$U_0 \stackrel{\text{def}}{=} \tilde{U}_0 - \Gamma^*(\Gamma\Gamma^*)^{-1}\Gamma\tilde{U}_0$$

corresponds to a function  $(u_0)_{\Omega_0} \in H_0^4(\Omega_0)$

- Choose  $u_0$  as the zero extension of the function  $(u_0)_{\Omega_0} \in H_0^4(\Omega_0)$ .
- Hence,  $u_0 \in H^4(\mathbb{R}^2) = H$ .

# Construction of an approximate inverse $\mathbb{A}$ for $D\mathbb{F}(u_0)$

- Approximate inverse  $\mathbb{A} : L_e \rightarrow H_e$  of  $D\mathbb{F}(u_0) = \mathbb{L} + D\mathbb{G}(u_0)$
- $\mathbb{A} = \mathbb{L}^{-1}\mathbb{B}$  where  $\mathbb{B} : L_e \rightarrow L_e$  approximates the inverse of  $I_d + D\mathbb{G}(u_0)\mathbb{L}^{-1}$
- Choose  $\mathbb{B} = I_d + \mathbb{B}_{\Omega_0}$  where  $\mathbb{B}_{\Omega_0} = \mathbf{1}_{\Omega_0} \mathbb{B}_{\Omega_0} \mathbf{1}_{\Omega_0} (\star)$

## Fourier series representation of $\mathbb{B}_{\Omega_0}$

If  $\mathbb{B}_{\Omega_0} : L_e \rightarrow L_e$  is a bounded linear operator such that  $(\star)$  holds, then  $\mathbb{B}_{\Omega_0}$  has a unique Fourier series representation  $B_p = (b_{i,j})_{i,j \in \mathbb{Z}^2} : \ell_e^2 \rightarrow \ell_e^2$  given by

$$b_{p,j} \stackrel{\text{def}}{=} \frac{1}{|\Omega_0|} \int_{\Omega_0} (\mathbb{B}_{\Omega_0} \psi_j)(x) \overline{\psi_i(x)} dx$$

where  $\psi_n(x) \stackrel{\text{def}}{=} e^{i\frac{\pi}{d}n \cdot x}$  for all  $n \in \mathbb{Z}^2$  and all  $x \in \Omega_0$ .

# Construction of an approximate inverse $\mathbb{A}$ for $D\mathbb{F}(u_0)$

- Numerically, we build  $B_p$  such that  $I_d + B_p \approx (I_d + D\mathbb{F}_p(U_0)L_p^{-1})^{-1}$
- $B_p$  has a representation  $\mathcal{F}^{-1}(B_p) = \mathbb{B}_{\Omega_0} : L_e \rightarrow L_e$  and  
 $\mathbb{B}_{\Omega_0} = \mathbf{1}_{\Omega_0} \mathbb{B}_{\Omega_0} \mathbf{1}_{\Omega_0}$
- Build  $\mathbb{B} = I_d + \mathbb{B}_{\Omega_0} : L_e \rightarrow L_e$
- Build  $\mathbb{A} = \mathbb{L}^{-1}\mathbb{B} : L_e \rightarrow H_e$
- Denote  $A_p = L_p^{-1}(I_d + B_p) \approx D\mathbb{F}_p(U_0)^{-1}$

# The bound $\mathcal{Y}_0$

- $\|\mathbb{A}\mathbb{F}(u_0)\|_l = \|\mathbb{L}\mathbb{A}\mathbb{F}(u_0)\|_2 = \|\mathbb{B}\mathbb{F}(u_0)\|_2$
- $\mathbb{F}(u_0)$  has a Fourier series representation  $F_p(U_0)$
- Using Parseval's identity,  
$$\|\mathbb{B}\mathbb{F}(u_0)\|_2^2 = |\Omega_0| \|(I_d + B_p)F_p(U_0)\|_2^2 = |\Omega_0| \|A_p F_p(U_0)\|_l^2$$
- Hence, the rigorous upper bound  $\mathcal{Y}_0$  is obtained via Fourier series and interval arithmetic computations.

# The bound $\mathcal{Z}_1$

We notice that

- $\|I_d - \mathbb{A}D\mathbb{F}(u_0)\|_l \leq \|I_d - A_p D F_p(U_0)\|_l + \mathcal{Z}_{1,1}$ .
- Note that  $\mathcal{Z}_{1,1} = O(e^{-ad})$ , which comes from the fact that

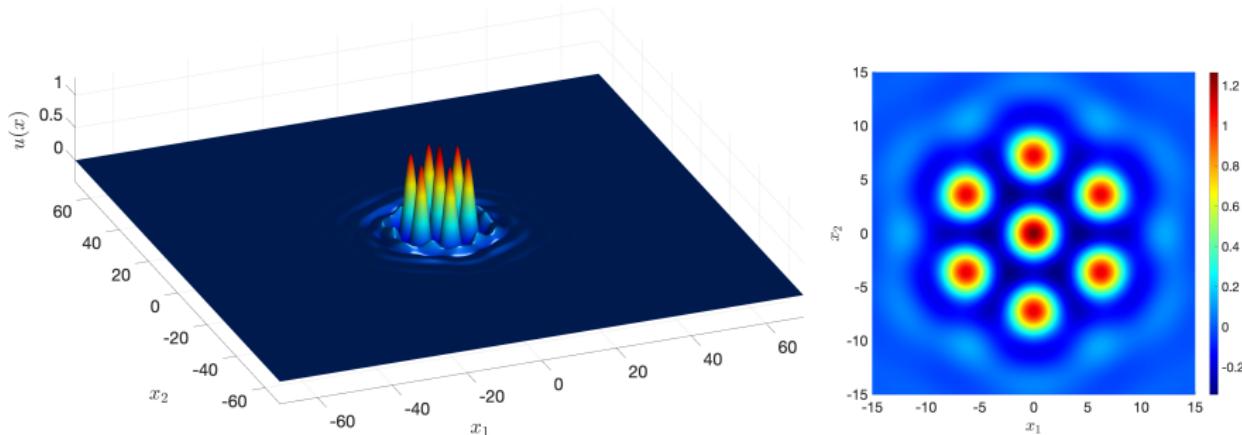
$$|\mathcal{F}^{-1}\left(\frac{1}{l}\right)(x)| \leq C e^{-a|x|}$$

where  $a$  is the smallest imaginary part of the roots of  $l$ .

- $\|I_d - A_p D F_p(U_0)\|_l \leq Z_1$  is computed with Fourier series computations.

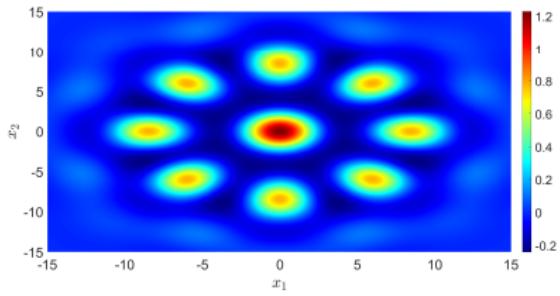
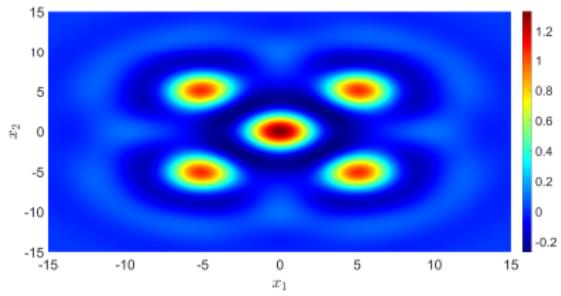
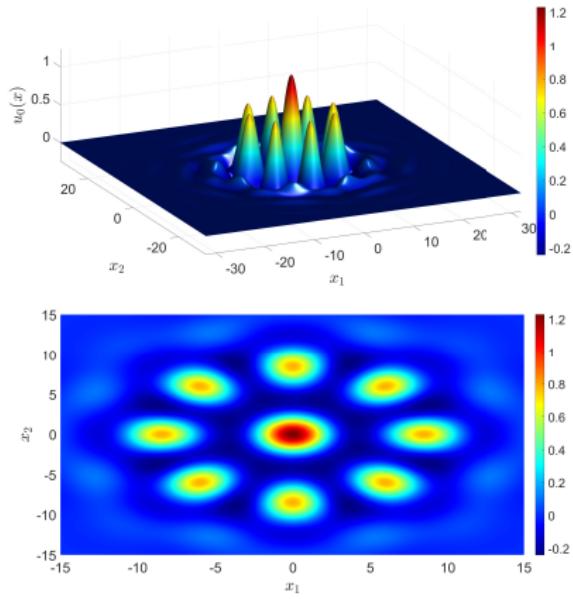
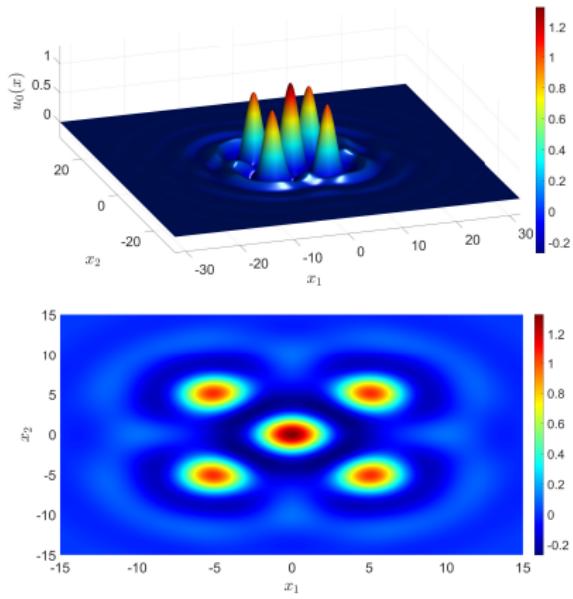
# Constructive proofs of existence of localized patterns in SH

$$(1 + \Delta)^2 u + \mu u + \nu_1 u^2 + \nu_2 u^3 = 0.$$



- Parameter values:  $\mu = 0.32$ ,  $\nu_1 = -1.6$  and  $\nu_2 = 1$
- Size of the Fourier projection:  $N = 130$ , Domain size:  $d = 70$  (vectors of size  $(N + 1)^2 = 17161$ )
- Bounds:  $\mathcal{Y}_0 = 2.7 \times 10^{-5}$ ,  $\mathcal{Z}_1 = 0.02$ ,  $\mathcal{Z}_2 = 212$
- $r = 2.8 \times 10^{-5}$  (rigorous global error bound in  $H^l$ )

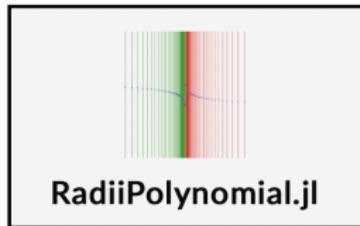
# Constructive proofs of existence of localized patterns in SH



# Rigorous interval arithmetic computations using Julia



O. Hénot  
McGill



# References

- ① Matthieu Cadiot, J.-P. L. and Jean-Christophe Nave. *Rigorous computation of solutions of semi-linear PDEs on unbounded domains via spectral methods*, arXiv:2302.12877.
- ② Dan J. Hill, Jason J. Bramburger and David J.B. Lloyd. *Approximate localised dihedral patterns near a Turing instability*. Nonlinearity, 36(5):2567, 2023.
- ③ Jan Bouwe van den Berg and J.-P. L. *Rigorous numerics in dynamics*. Notices Amer. Math. Soc. 62 (2015), no. 9, 1057–1061.
- ④ Javier Gómez-Serrano. *Computer-assisted proofs in PDE: a survey*. SeMA J. 76 (2019), no. 3, 459–484.



**NSERC  
CRSNG**



CENTRE  
DE RECHERCHES  
MATHÉMATIQUES